

## Ideal Gas Law:

- Below  $\sim 100$  km, to a good approximation the atmosphere behaves as ideal gas
- Mean free path:  $\lambda \sim 10^{-7}$  m @  $z=0$  km,  $\lambda \sim 1$  m @  $z=100$  km

So:  $pV = nR^*T$

$p$  = pressure,  $V$  = volume,  $n$  = no. of moles  
 $R^*$  = universal gas constant =  $8.31 \text{ J kg}^{-1} \text{ } ^\circ\text{C}^{-1}$   
 $T$  = absolute temperature

- For the atmosphere, it is easier to deal with densities, so gas law is gas specific:

$$R = \frac{R^*}{M_d}$$

$$M_d = \text{molecular weight of dry air}$$

$$\approx 80\% \times 28 + 20\% \times 32$$

$$\begin{array}{ccc} \uparrow & & \uparrow \\ \text{N}_2 & & \text{O}_2 \end{array}$$

$$\approx 29 \text{ g/mol}$$

$$\Rightarrow R = \frac{8.31}{29/1000}$$

$$\approx 287 \text{ J kg}^{-1} \text{ } ^\circ\text{C}^{-1}$$

$$\Rightarrow pV = \underbrace{n \times M_d}_{\text{mass}} \times R \times T$$

$$\Rightarrow \boxed{p = \rho RT}$$

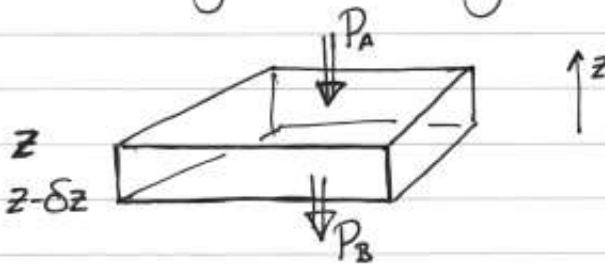
↑  
density

Equation of State for  
Atmosphere

Also, sometimes see:  $\alpha = \frac{1}{\rho} = \text{"specific volume"}$

Hydrostatic Eq.:

- ideal gas law immediately allows us to see how pressure varies with altitude



imagine slab of the atmosphere. The force on face B equals force on face A plus the weight of the slab

$$\Rightarrow p_B \times A = p_A \times A + g \cdot \rho \cdot A \cdot \delta z \quad \textcircled{1}$$

A = area of face A, B

Since  $p = p(z)$ , can write

$$p_0 = p_1 - \frac{dp}{dz} \delta z \quad \left(1^{\text{st}} \text{ order Taylor expansion}\right)$$

① becomes:

$$\Rightarrow \left(p_1 - \frac{dp}{dz} \delta z\right) A = p_1 \cdot A + A g \rho \delta z$$

$$\Rightarrow \boxed{\frac{dp}{dz} = -g\rho} \quad \text{Hydrostatic equation}$$

→ Using ideal gas law  $p = \rho RT$

$$\frac{dp}{dz} = -\frac{g\rho}{RT}$$

$$\Rightarrow \frac{dp}{p} = -\frac{g}{RT} dz$$

Note  $\frac{RT}{g}$  has dimensions [m]

Let  $H = \frac{RT}{g} = \text{ATMOSPHERIC SCALE HEIGHT} \sim \text{constant (why)}$

$$\sim \frac{287 \times 300}{9.8}$$

$\sim 9 \text{ km}$  in tropics (7.5 km at poles)

So  $\frac{dp}{p} \approx -\frac{dz}{H}$   $H$  const.

integrate from surface  $z=0$ ,  $p=p_s$  to altitude  $z$ , pressure  $p$

$$\int_{p_s}^p \frac{dp}{p} = - \int_0^z \frac{dz}{H}$$

$$\ln p/p_s = -z/H$$

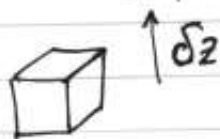
$$\Rightarrow \boxed{p = p_s e^{-z/H}}$$

So pressure decays exponentially away from the surface with a characteristic scale,  $H$ .

Geopotential ~~is~~ - turns out to be a useful way of representing height

$\phi$  = work done against gravity in lifting 1 kg parcel of air from surface to altitude  $z$ .

= geopotential



work done = ~~mass~~ force  $\times$  distance

$$\phi = \int_0^z m g dz$$

n.b.  $g$  not constant  
why? 2 reasons

$$\phi = \int_{p_s}^p \frac{dp}{\rho}$$

(using ~~the ideal gas law~~ hydrostatic equation)

n.b. geopotential depends only on the elevation the parcel is at, and ~~not~~ on how it got there.

Can define geopotential height:

$$z = \phi(z)/g_0$$

$$g_0 = 9.8 \text{ ms}^{-2}$$

Value of geopotential height:

back to hydrostatic eq<sup>n</sup>:

$$\frac{dp}{dz} = -g\rho$$

from def<sup>n</sup> of geopotential  $d\phi = g dz$

n.b. looks a lot like hydrostatic eq<sup>n</sup>, but have got rid of density

$$\Rightarrow d\phi = -RT \frac{dp}{p}$$

$$\Rightarrow \phi_2 - \phi_1 = -R \int_{p_1}^{p_2} T \frac{dp}{p}$$

$$z_2 - z_1 = \frac{R\bar{T}}{g_0} \ln p_1/p_2 \quad \bar{T} = \text{average temperature}$$

$\Rightarrow$  The mean temperature in a layer of the atmosphere is proportional to the GEOPOTENTIAL <sup>HEIGHT</sup> THICKNESS

Example:

What is the geopotential height thickness between 1000 mb & 500 mb which would mean the average temperature in that layer was freezing?

$$\Delta z = \frac{287 \times 273}{9.8} \ln \frac{1000}{500}$$

$$\approx 5500 \text{ m}$$

Forecasters use the 1000 mb-500 mb thickness as an indicator for ~~detecting~~ determining the rain-snow line in winter...